

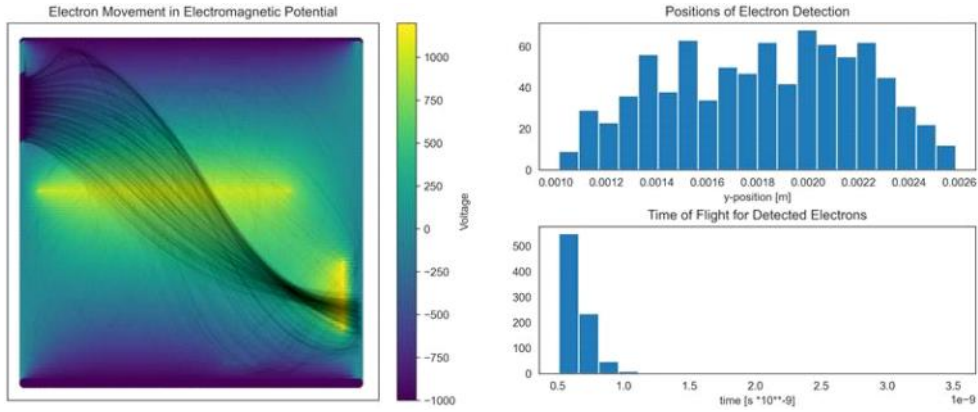
1-D Hydrodynamics

Monday, 4 December 2023 10:56

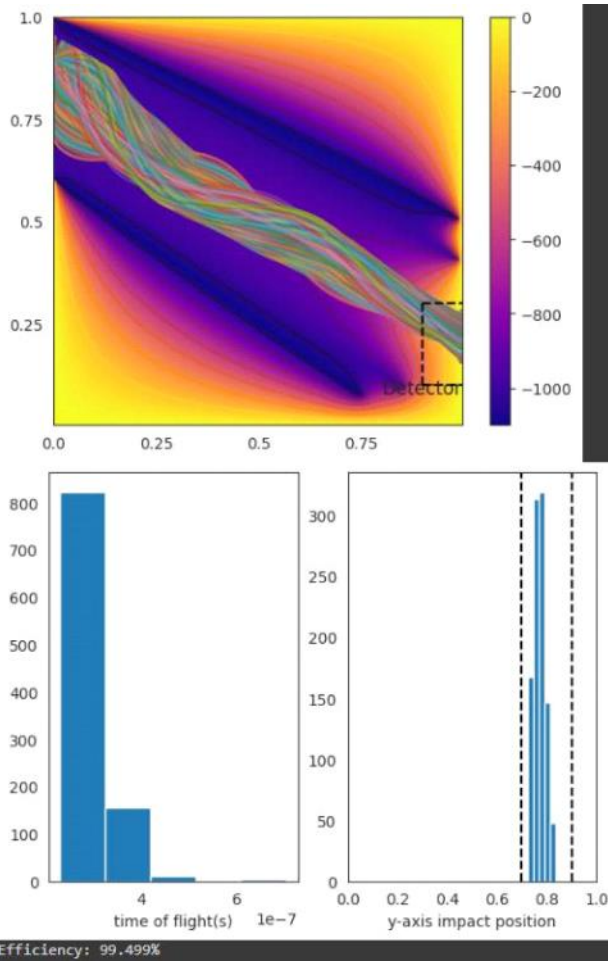
Design Prize:

Honorable Mentions:

Rafaela Lambrinos Maria: for a functional and interesting design!

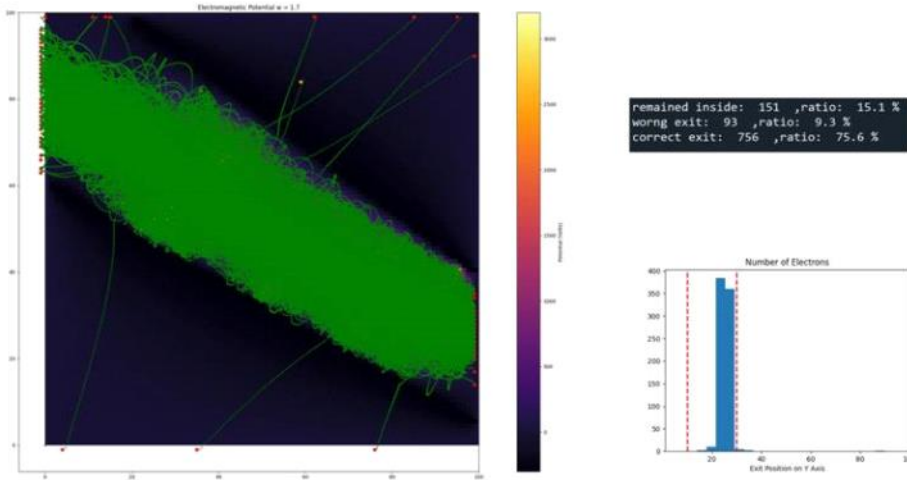


Luigi Willi Pietro: Bi-cubic interpolation! (attempt) Very high detection rate with a funnel-like design.



Winner:

Timo Elias Rieger: Best documented design, with plenty of explorations!



Recall: LAX Method

$$\rho_i^{(n+1)} = \frac{1}{2}(\rho_{i+1}^{(n)} + \rho_{i-1}^{(n)}) - \frac{c}{2}(\rho_{i+1}^{(n)} - \rho_{i-1}^{(n)})$$

$$c = \frac{\Delta t u}{\Delta x}$$

For linear advection $f(\rho) = \rho u$

Rewrite the above in terms of fluxes

$$\boxed{A} \quad \rho_i^{(n+1)} = \frac{1}{2}(\rho_{i+1}^{(n)} + \rho_{i-1}^{(n)}) - \frac{\Delta t}{2\Delta x}(F_{i+1/2}^{(n)} - F_{i-1/2}^{(n)})$$

also Rewrite the 1st order upwind method:

$$\boxed{A'} \quad \rho_i^{(n+1)} = \rho_i^{(n)} - \frac{\Delta t}{\Delta x}(F_i^{(n)} - F_{i-1}^{(n)}) \text{ for } u \geq 0$$

$$- \frac{\Delta t}{\Delta x}(F_{i+1}^{(n)} - F_i^{(n)}) \text{ for } u < 0$$

Now for 1-D hydrodynamics, for 3 conserved quantities:

1. Cons. of Mass: $F = \rho u$

2. Cons. of Momentum: $F = \rho u^2 + P$ - pressure

3. Cons. of Energy: $F = u(E + P)$

Note u doesn't have to be a constant here.

Consider a state vector

$$\underline{U} = \begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix} \quad F(U) = \begin{bmatrix} \rho u \\ \rho u^2 + P \\ u(E + P) \end{bmatrix}$$

----- | non diagonal

$$\frac{\partial u}{\partial t} + \nabla \cdot F(u) = 0$$

Very general equation of "stuff"

↳ 3 equations

(in our 1-D case we have $\nabla \cdot \equiv \frac{\partial}{\partial x}$)

Unknowns: $\rho, \rho u, E, P$

4 unknowns ...

Equation of State

$E = \text{Kinetic Energy} + \text{Thermal Energy}$

$$\frac{1}{2} \rho u^2 + \rho e$$

thermal energy density per unit mass (specific)

$\rho, \rho u, e, P$

$$PV = nRT \quad (\text{Ideal Gas})$$

(same)

$$e = \frac{P}{(\gamma-1)\rho}$$

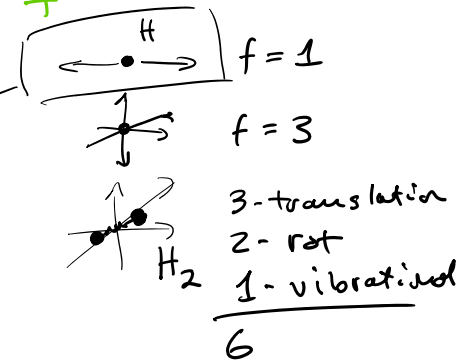
$\gamma = \frac{f+2}{f}$ degrees of freedom.

$\gamma = 3$

$$e = \frac{1}{2} \frac{P}{\rho}$$

$$E = \frac{1}{2} \rho u^2 + \frac{1}{2} P$$

$$F(u) = \left[\rho u, \rho u^2 + P, u \left(\frac{1}{2} \rho u^2 + \frac{3}{2} P \right) \right]$$



Think about Finite Volumes:

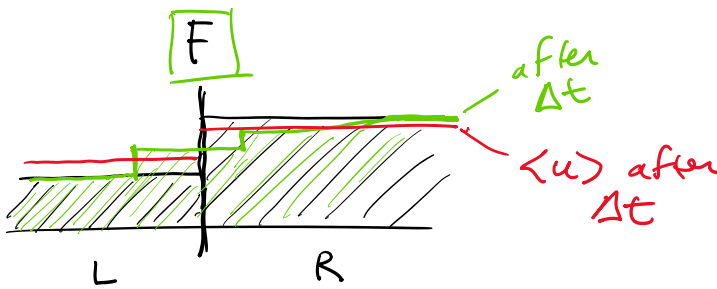
u is representing the average value of the "material" in the cell:

$$\langle u \rangle_i^n = \frac{1}{\Delta x} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} u(x, t^n) dx$$

$$\langle u \rangle_i^n = \frac{1}{\Delta x} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} u(x, t^n) dx$$

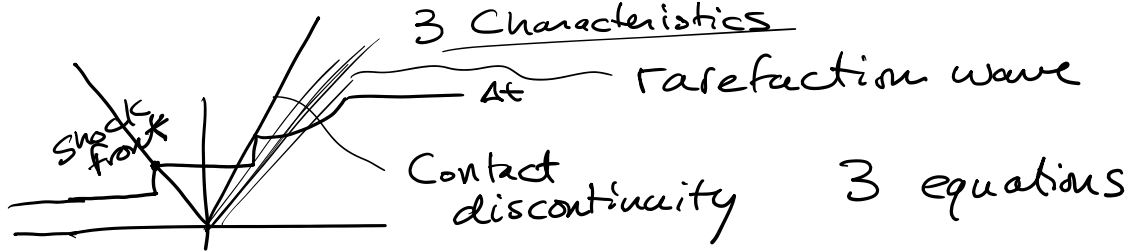
in actual fact we should consider this average over a timestep.

$$\frac{\langle u \rangle_i^{n+1} - \langle u \rangle_i^n}{\Delta t} + \frac{F_{i+\frac{1}{2}}^n - F_{i-\frac{1}{2}}^n}{\Delta x} = 0$$

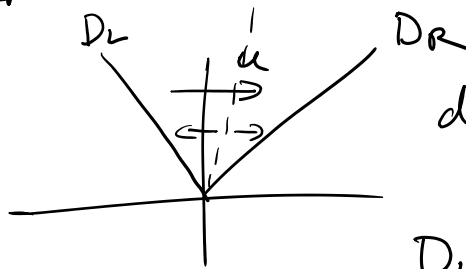


$F(u_L, u_R)$
Riemann Solver

★ Good Approximation



Approximate Riemann Solver :



discontinuity speeds of the left and right waves.

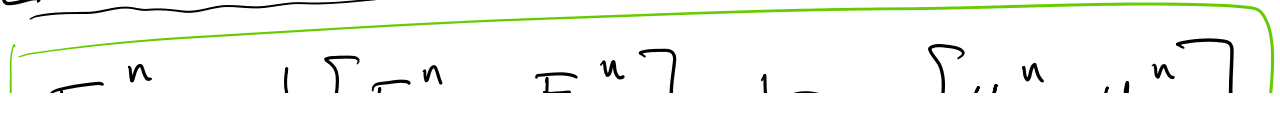
$$D_L = u - c_s$$

$$D_R = u + c_s$$

$$D_{max} = \max(D_L, D_R) = |u| + c_s$$

$$c_s = \sqrt{\gamma \frac{P}{\rho}} \text{ speed of sound}$$

Lax-Friedrichs Riemann Solver



$$F_{i-\frac{1}{2}}^n = \frac{1}{2} [F_i^n + F_{i-1}^n] - \frac{1}{2} D_{\max} [u_i^n - u_{i-1}^n]$$

Method B

$$\langle u \rangle_i^{n+1} = \langle u \rangle_i^n - \frac{\Delta t}{\Delta x} [F_{i+\frac{1}{2}}^{n+\frac{1}{2}} - F_{i-\frac{1}{2}}^{n+\frac{1}{2}}]$$

Method C (dop <>)

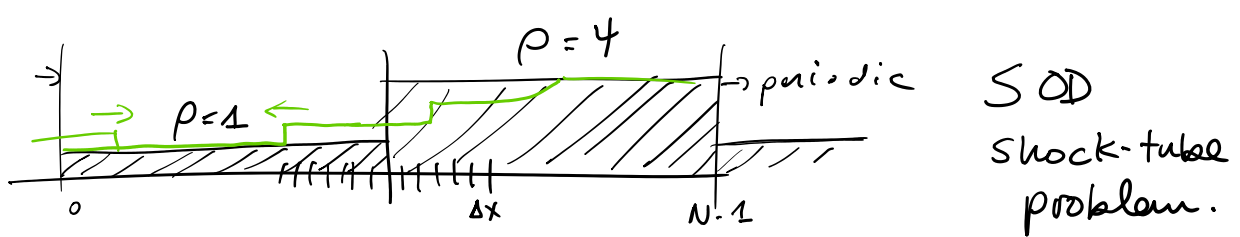
$$*u_i^{n+\frac{1}{2}} = *u_i^n - \frac{\Delta t}{2\Delta x} [F_{i+\frac{1}{2}}^n - F_{i-\frac{1}{2}}^n]$$

"predictor step"

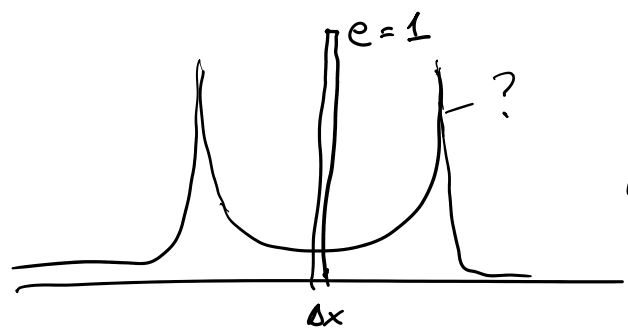
$$u_i^{n+1} = u_i^n - \frac{\Delta t}{\Delta x} [F(*u_i, *u_{i+1}) - F(*u_{i-1}, *u_i)]$$

"corrector step"

use
LAX-Friedrichs
Riemann Solver



$u=0$
 $\epsilon = 10^{-5}$ — solve for P given p
 $\Delta t ?$ $D_{\max} < \frac{\Delta x}{\Delta t}$



Sedov-Taylor Blast wave
 (piston or bomb)

Sedov-Taylor Blast wave
solution (like a bomb)